Numerical Simulation of Counterstreaming Plasma Interactions using a Multifluid Model

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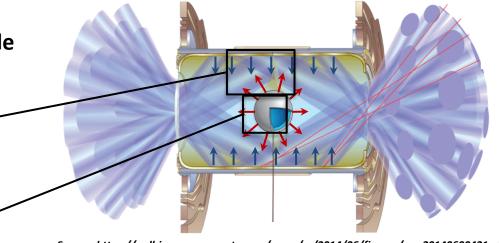
Background and Motivation

Inertial Confinement Fusion: Colliding plasmas from hohlraum wall and capsule

Interpenetration of plasma flows from capsule and hohlraum wall

- Large range of Z: $2 \le Z \le 60$
- Supersonic flows ($\Delta u \approx 10^8$ cm/s)

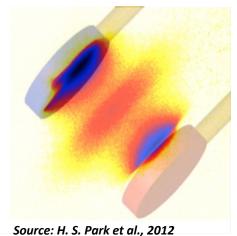
Species separation inside target capsule



Source: https://csdl-images.computer.org/mags/cs/2014/06/figures/mcs20140600421.gif

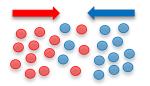
High Energy Density Physics (HEDP) Experiments

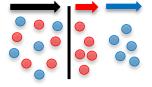
Carbon plasma streams ablating off paddles hit by laser beams and colliding with each other



Interpenetrating plasmas

Multifluid phenomena that we want to model





Plasma species separation





Current simulation tools are not sufficiently versatile

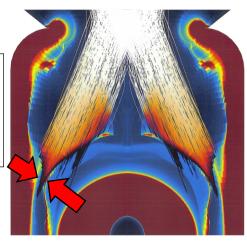
Lack key physics

Single-Fluid Multi-species Hydrodynamic Solvers

- o Example: HYDRA, LASNEX
- Single velocity field insufficient to model multiple inter-penetrating fluids
- Unphysical shocks

Simulation of plasma dynamics in hohlraum using HYDRA

Density pile—up predicted when plasma streams collide



Current workarounds: species diffusion models

Collisional Kinetic Solvers

- Example: LOKI, OSIRIS, PSC
- High computational cost to simulate small volumes
- Impractical for experimental scales



Governing Equations: We solve the inviscid Euler equations for each ion species

$$\frac{\partial \rho_{\alpha}}{\partial t} + \nabla \cdot (\rho_{\alpha} \mathbf{u}_{\alpha}) = 0$$

$$\frac{\partial \rho_{\alpha} \mathbf{u}_{\alpha}}{\partial t} + \nabla \left(P_{\alpha} + \rho_{\alpha} \mathbf{u}_{\alpha} \otimes \mathbf{u}_{\alpha}\right) = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \nabla \phi + \sum_{\beta \neq \alpha} \mathbf{R}_{\alpha,\beta} & \text{Interaction} \\ between species \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] = \begin{bmatrix} -Z_{\alpha} e n_{\alpha} \mathbf{u}_{\alpha} \cdot \nabla \phi + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta}\right) \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left[\left(\mathcal{E}_{\alpha} + P_{\alpha}\right) \mathbf{u}_{\alpha} \right] \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} +$$

Assuming quasineutral, isothermal electrons*

$$\nabla \phi = \frac{T_e}{n_e} \nabla n_e + \frac{1}{n_e} \sum_{\alpha} R_{e,\alpha}$$

Electron momentum equation neglecting inertia terms and assuming

$$P_e = n_e T_e$$

Frictional drag

$$\mathbf{R}_{\alpha,\beta} = m_{\alpha} n_{\alpha} \nu_{\alpha,\beta} \left(\mathbf{u}_{\beta} - \mathbf{u}_{\alpha} \right)$$

Frictional heating and thermal equilibration

$$egin{aligned} Q_{lpha,eta} &= Q_{lpha,eta}^{
m fric} + Q_{lpha,eta}^{
m eq} \ Q_{lpha,eta}^{
m fric} &= m_{lpha,eta} n_{lpha}
u_{lpha,eta} \left(\mathbf{u}_{eta} - \mathbf{u}_{lpha}
ight)^2 \ Q_{lpha,eta}^{
m eq} &= -3 m_{lpha} n_{lpha} rac{
u_{lpha,eta}}{m_{lpha} + m_{eta}} \left(T_{lpha} - T_{eta}
ight) \end{aligned}$$





Reformulated Governing Equations

Ion Euler equations with isothermal, quasineutral e

Advective nature of electrostatic force



- o Included electron pressure on LHS with hydrodynamic pressure
- o Derived the eigenstructure for characteristic-based discretization

Effect of discretization error in dense species on **dynamics of sparse species**



Reformulation of electrostatic source terms to avoid sums/differences of terms of disparate scales

$$\begin{split} \frac{\partial \rho_{\alpha}}{\partial t} + \nabla \cdot (\rho_{\alpha} \mathbf{u}_{\alpha}) &= 0, \\ \frac{\partial \rho_{\alpha} \mathbf{u}_{\alpha}}{\partial t} + \nabla \left(\rho_{\alpha} \mathbf{u}_{\alpha} \otimes \mathbf{u}_{\alpha} + \frac{\mathbf{P}_{\alpha}^{*}}{\mathbf{P}_{\alpha}^{*}} \right) &= Z_{\alpha} T_{e} n_{e} \nabla \left(\frac{n_{\alpha}}{n_{e}} \right) + \frac{Z_{\alpha} n_{\alpha}}{n_{e}} \sum_{\beta} \mathbf{R}_{e,\beta} + \mathbf{R}_{\alpha,e} + \sum_{\beta \neq \alpha} \mathbf{R}_{\alpha,\beta}, \\ \frac{\partial \mathcal{E}_{\alpha}}{\partial t} + \nabla \cdot \left\{ \left(\mathcal{E}_{\alpha} + \frac{\mathbf{P}_{\alpha}^{*}}{\alpha} \right) \mathbf{u}_{\alpha} \right\} &= Z_{\alpha} T_{e} n_{e} \nabla \left(\frac{\mathbf{u}_{\alpha} n_{\alpha}}{n_{e}} \right) + \frac{Z_{\alpha} n_{\alpha}}{n_{e}} \sum_{\beta} \mathbf{u}_{\alpha} \cdot \mathbf{R}_{e,\beta} + \sum_{\beta \neq \alpha} \left(\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,\beta} \right) \\ &+ \mathbf{R}_{\alpha,e} \cdot \mathbf{u}_{\alpha} + Q_{\alpha,e}^{eq}, \end{split}$$

where $P_{\alpha}^{*}=P_{\alpha}+Z_{\alpha}T_{e}n_{\alpha}$ [Electron pressure] is the "augmented pressure" (hydro + e^{-})

Wavespeeds (eigenvalues) : $\mathbf{v}, \mathbf{v} \pm \sqrt{\frac{\gamma_{\alpha} P_{\alpha}^*}{\rho_{\alpha}}}$

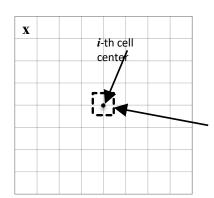




Summary of Numerical Method

High-Order Conservative Finite-Difference/Finite-Volume Method

4th order finite-volume discretization (using the CHOMBO library) with AMR



3D Domain
$$\Omega \equiv \{\mathbf{x}: 0 \leq \mathbf{x} \cdot \mathbf{e}_d \leq L_d, 1 \leq d \leq 3\}$$

discretized into computational cells

$$\mathbf{-} \omega_{\mathbf{i}} = \prod_{d=1}^{3} \left[\left(\mathbf{i} - \frac{1}{2} \mathbf{e}_{d} \right) h, \left(\mathbf{i} + \frac{1}{2} \mathbf{e}_{d} \right) h \right]$$

i: 3-dimensional integer index (i, j, k) h: grid spacing

Spatially-discretized ODE in time (integrated in time using 4th order Runge-Kutta method)

$$\frac{\partial \bar{\mathbf{u}_i}}{\partial t} = \frac{1}{h} \sum_{d=1}^{3} \left(\left\langle \hat{\mathbf{F}_{i+\frac{1}{2}\mathbf{e}_d}} \right\rangle - \left\langle \hat{\mathbf{F}_{i-\frac{1}{2}\mathbf{e}_d}} \right\rangle \right)$$
 Cell-averaged solution Face-averaged fluxes

Strong shocks and gradients
O(1) to O(1e-14)



- Characteristic-based discretization
- Implemented 5th-order WENO scheme with Monotonicity-Preserving limiting



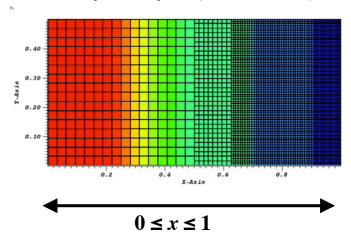
Example: Collisionless Electrostatic Single Fluid Shock Tube

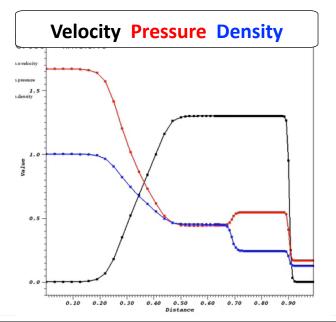
Extension of the Sod's shock tube test case

- o Initial solution: Riemann problem
- Helium gas specified
- Essentially a 1D problem
- Inviscid wall along x, periodic along y and z.

Since **Debye length much smaller than domain**, dynamics *similar to neutral gas dynamics*with *hydro pressure augmented by the electron pressure*

Density color plot (0.125 to 1.0)





Initial Riemann discontinuity decomposes into a shock, a contact, and a rarefaction





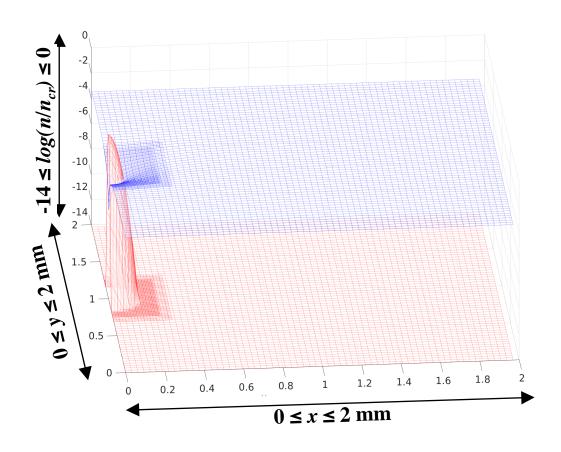
Example: Single Species Expansion into Gas Fill with AMR – *Initial Setup*

Expansion of a carbon blob in the presence of helium gas fill (2D)

- Initial solution: a carbon species piled up on one end (Gaussian blob density); gas fill present in the space everywhere else.
- Boundary conditions: Solid wall
 BCs along x and y

Reference quantities:

Mass: *proton mass* (1.6730e-24 g); Number density: n_{crit} (9.0320e+21 cm⁻³); Length: 1 mm; Time: 3.2314e-09 s; Temperature: 1 keV (1.6022e-09 ergs)









Example: Single Species Expansion into Gas Fill with AMR – Solution Evolution

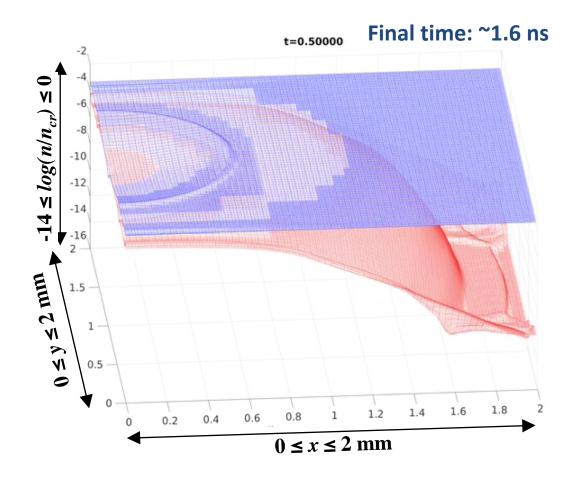
Expansion of a carbon blob in the presence of helium gas fill (2D)

- Initial solution: a carbon species piled up on one end (Gaussian blob density); gas fill present in the space everywhere else.
- Boundary conditions: Solid wall
 BCs along x and y

AMR: Refined mesh adaptively generated in regions of high gradients

Reference quantities:

Mass: *proton mass* (1.6730e-24 g); Number density: n_{crit} (9.0320e+21 cm⁻³); Length: 1 mm; Time: 3.2314e-09 s; Temperature: 1 keV (1.6022e-09 ergs)



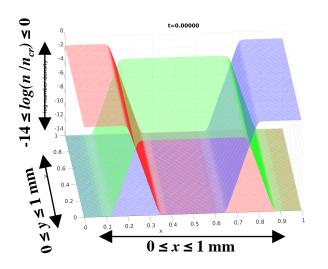




Example: Two Species Interpenetration with Gas Fill Problem Setup

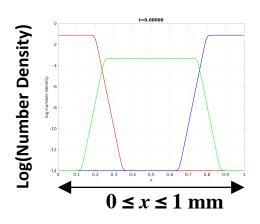
Interpenetration of carbon and carbon streams in the presence of helium gas fill (2D)

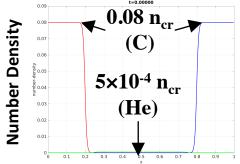
- o **Initial solution:** two species piled up on either end (*smoothed slab* density); gas fill present in the space in between.
- \circ **Temperature variation along y** the plasmas are hotter in the center of the domain

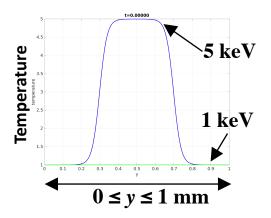


Boundary conditions:

- Solid wall BCs along x
- Periodic along y







Reference quantities:

Mass: proton mass (1.6730e-24 g)

Number density: n_{crit} (9.0320e+21 cm⁻³)

Length: 1 mm

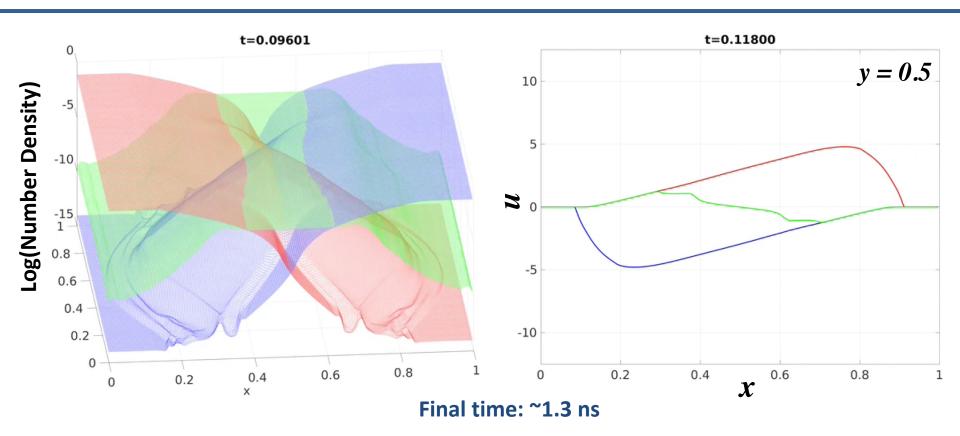
Temperature: 1 keV (1.6022e-09 ergs)







Example: Two Species Interpenetration with Gas Fill



- Species interaction prevents one species from reaching the other end of the domain along x
- The fill gas is pushed towards the center of the domain by the carbon streams.

Reference quantities:

Number density: n_{crit} (9.0320e+21 cm⁻³);

Length: 1 mm; Time: 3.2314e-09 s;

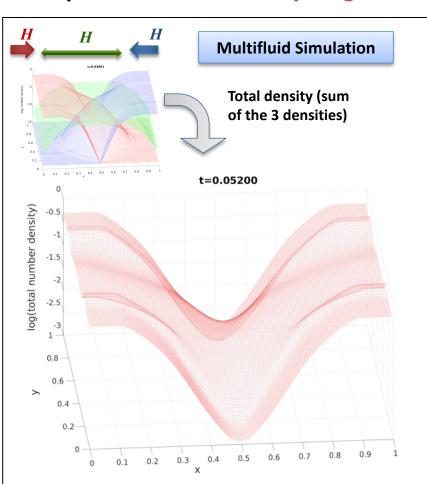
Velocity: 3.0946e+07 cm/s

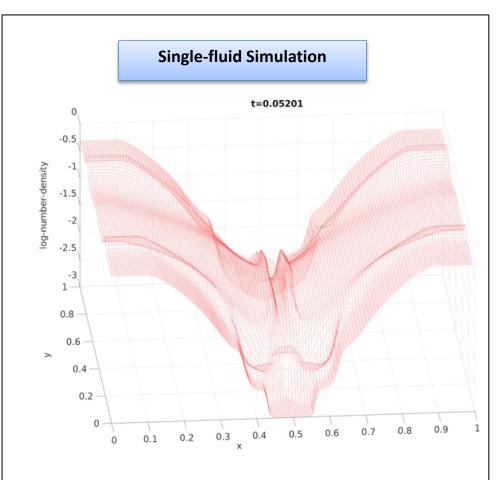




Multifluid vs. Single Fluid Simulations - How do the solutions differ?

Interpenetration of two hydrogen streams in the presence of hydrogen gas fill (2D)









Conclusions and Future Work

Summary

EUCLID: EUlerian Code for pLasma Interpenetration Dynamics

- Developed a 3D, parallel, AMR-capable multifluid flow solver
- o Implemented the *quasineutral, isothermal electron model* as a computationally tractable electron model for our target applications.
- Verified EUCLID for accuracy and convergence (benchmark cases, manufactured solutions)
- Simulated flows motivated by laboratory astrophysics experiments and ICF hohlraums.

Current and Future Work

- Conduct simulations of plasma interpenetration experiments (e.g. Ross et al., 2013, Le Pape et al., ongoing)
- Investigate the use of IMEX time integrators for stiff collisional terms involving high-Z species.
- o Investigate *higher-fidelity electron models*, for example, adding an electron energy equation.
- Add source terms to energy equations to simulate heating







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